Search for Top Squarks in *R*-Parity-Violating Supersymmetry Using Three or More Leptons and *b*-Tagged Jets

S. Chatrchyan *et al.**
(CMS Collaboration)

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A search for anomalous production of events with three or more isolated leptons and bottom-quark jets produced in pp collisions at $\sqrt{s} = 8$ TeV is presented. The analysis is based on a data sample corresponding to an integrated luminosity of 19.5 fb^{-1} collected by the CMS experiment at the LHC in 2012. No excess above the standard model expectations is observed. The results are interpreted in the context of supersymmetric models with signatures that have low missing transverse energy arising from light top-squark pair production with R-parity-violating decays of the lightest supersymmetric particle. In two models with different R-parity-violating couplings, top squarks are excluded below masses of 1020 GeV and 820 GeV when the lightest supersymmetric particle has a mass of 200 GeV.

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Supersymmetric (SUSY) extensions of the standard model (SM) solve the hierarchy problem while unifying particle interactions [1,2]. Among SUSY models, "natural" supersymmetry refers to those characterized by small fine-tuning needed to describe particle spectra. It requires top squarks (stops), the top-quark superpartners, to be lighter than about 1 TeV. These models have received substantial interest in light of the discovery of a Higgs boson with mass near 125 GeV [3,4] because the stop should be the superpartner most strongly coupled to the Higgs boson.

Natural models feature pair production of stops that decay to a number of final states. To fully test supersymmetric naturalness, searches for all possible decay chains should be carried out. These can be broadly categorized as R-parity conserving (RPC) or violating (RPV) [5], where R-parity is defined by $R = (-1)^{3B+L+2s}$, with B and L the baryon and lepton numbers, and s the particle spin. All SM particle fields have $R = +1$ while all superpartner fields have $R = -1$. When R-parity is conserved, superpartners are produced in pairs, the lightest superpartner (LSP) is stable and a dark-matter candidate, and proton stability is ensured. Most recent searches for naturalness have focused on RPC models [6–8].

Supersymmetric models with RPV interactions violate either B or L but can avoid proton decay limits [9,10]. The superpotential W_{RPV} includes three trilinear terms parametrized by the Yukawa couplings λ_{ijk}, λ'_{ijk}, and λ''_{ijk}:

$$W_{\mathrm{RPV}} = \frac{1}{2}\lambda_{ijk}L_iL_j\bar{E}_k + \lambda'_{ijk}L_iQ_j\bar{D}_k + \frac{1}{2}\lambda''_{ijk}\bar{U}_i\bar{D}_j\bar{D}_k, \quad (1)$$

where i, j, and k are generation indices; L and Q are the $SU(2)_L$ doublet superfields of the lepton and quark; and the $\bar{E}$, $\bar{D}$, and $\bar{U}$ are the $SU(2)_L$ singlet superfields of the charged lepton, downlike quark, and uplike quark. The third term violates baryon number conservation, while the first two terms violate lepton number conservation. These terms do not preclude a natural hierarchy [11].

The RPV interactions allow for single production of SUSY particles (sparticles) and for sparticle decay into SM only particles. The latter is explored in this Letter. Prior searches for RPV interactions in multilepton final states include those at LEP [12–14], the Tevatron [15,16], at HERA [17,18], and at the Large Hadron Collider (LHC) [19–21].

Because the LSP is unstable in RPV models, a common experimental strategy of SUSY searches—selecting events with large missing transverse energy ($E_{\mathrm{T}}^{\mathrm{miss}}$)—is not effective [9]. Instead, we use S_{T}, the scalar sum of $E_{\mathrm{T}}^{\mathrm{miss}}$ and the transverse energy of jets and charged leptons, to differentiate between signal and standard model backgrounds.

In this Letter we present the result of a search for pair production of top squarks with RPV decays of the lightest sparticle, using multilepton events and bottom-tagged (b-tagged) jets. The data set used here corresponds to an integrated luminosity of 19.5 fb^{-1}, recorded in 2012 with the CMS detector at the LHC in proton-proton collisions at a center-of-mass energy of 8 TeV.

The coordinate system in CMS is right handed, with the origin at the nominal interaction point. Pseudorapidity is given by $\eta \equiv -\ln[\tan(\theta/2)]$, where the polar angle θ is defined with respect to the counterclockwise beam direction. The azimuthal angle ϕ is measured relative to the direction to the center of the LHC ring.

The CMS detector [22] has cylindrical symmetry around the pp beam axis with tracking and muon detectors covering the pseudorapidity range $|\eta| < 2.4$. The tracking system measures the trajectory and momentum of charged

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*Full author list given at the end of the article.

particles and consists of multilayered silicon pixel and strip detectors in a 3.8 T solenoidal magnetic field. Particle energies are measured with concentric electromagnetic and hadron calorimeters, which cover $|\eta| < 3.0$ and $|\eta| < 5.0$, respectively. Muon detectors consisting of wire chambers are embedded in the steel return yoke outside the solenoid. The trigger thresholds in a two-level trigger system are tuned to accept a few hundred data events per second from the pp interactions.

We select events with three or more leptons (including tau leptons) that are accepted by a trigger requiring two light leptons, which may be electrons or muons. Any opposite-sign, same-flavor pair of electrons or muons must have an invariant mass $m_{\ell\ell} > 12$ GeV, removing low-mass bound states and $\gamma^* \to \ell^+\ell^-$ production.

Electrons and muons are reconstructed using the tracker, calorimeter, and muon systems. Details of reconstruction and identification can be found in Ref. [23] for electrons and in Ref. [24] for muons. We require that at least one electron or muon in each event have transverse momentum of $p_T > 20$ GeV. Additional electrons and muons must have $p_T > 10$ GeV and all of them must be within $|\eta| < 2.4$.

The majority of hadronic decays of tau leptons (τ_h) yield either a single charged track (one-prong) or three charged tracks (three-prong), occasionally with additional electromagnetic energy from neutral pion decays. We use one- and three-prong τ_h candidates that have $p_T > 20$ GeV, reconstructed with the "hadron plus strips" method [25]. Leptonically decaying taus are included with other electrons and muons.

To ensure that electrons, muons, and τ_h candidates are isolated, we use a particle-flow algorithm [26,27] to identify the source of transverse energy deposits in the trackers and calorimeters. We then sum the energy deposits in a cone of radius 0.3 in $\Delta R = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2}$ around the candidate and subtract the lepton p_T to calculate $E_{T,\mathrm{cone}}$. We remove energy from additional proton-proton collisions that occur simultaneously by subtracting a per-event correction [23,28]. For electrons and muons, we divide $E_{T,\mathrm{cone}}$ by the lepton p_T to find the relative isolation $I_{\mathrm{rel}} = E_{T,\mathrm{cone}}/p_T$, which has to be less than 0.15. We require $E_{\mathrm{cone}} < 2$ GeV for τ_h candidates.

We use jets reconstructed from particle-flow candidates [28] using the anti-k_T algorithm [29] with a distance parameter of 0.5, that have $|\eta| < 2.5$ and $p_T > 30$ GeV. Jets are required to be a distance $\Delta R > 0.3$ away from any isolated electron, muon, or τ_h candidate. To determine if the jet originated from a bottom quark, we use the combined secondary-vertex algorithm, which calculates a likelihood discriminant using the track impact parameter and secondary-vertex information. This discrimination selects heavy-flavor jets with an efficiency of 70% and suppresses light-flavor jets with a misidentification probability of 1.5% [30].

Monte Carlo (MC) simulations are used to estimate some of the SM backgrounds and to understand the efficiency and acceptance of the signal models. The SM background samples are generated using MADGRAPH [31] with parton showering and fragmentation modeled using PYTHIA (version 6.420) [32] and passed through a GEANT4-based [33] representation of the CMS detector. Signal samples [11] are generated with MADGRAPH and PYTHIA and passed through the CMS fast-simulation package [34]. Next-to-leading and next-to-leading-log cross sections and their uncertainties for the SUSY signal processes are from the LHC SUSY cross sections working group [35–39].

Multilepton signals have two main sources of backgrounds, the first arising from processes that produce genuine multilepton events. The most significant examples are WZ and ZZ production, but rare processes such as $t\bar{t}W^\pm$ and $t\bar{t}Z$ also contribute. We assess the contribution from these processes using samples simulated by MADGRAPH. Samples simulating WZ and ZZ have been validated in control regions in data. For the rarer background processes, we rely solely on simulation.

The second source originates from objects that are misclassified as prompt, isolated leptons, but are actually hadrons, leptons from a hadron decay, etc. Misidentified leptons are classified in three categories: misidentified light leptons (electrons and muons), misidentified τ_h leptons, and light leptons originating from asymmetric internal conversions. The methods used in this paper are described in more detail in Ref. [20]

We estimate the contribution of misidentified light leptons by measuring the number of isolated tracks and applying a scale factor between isolated leptons and isolated tracks. These scale factors are measured in control regions that contain leptonically decaying Z-bosons and a third, isolated track, as well as in control regions with opposite-sign, opposite-flavor leptons, which are $t\bar{t}$ dominated. The scale factor is then the probability for the third track to pass the lepton identification criteria. We find the scale factors to be $(0.9 \pm 0.2)\%$ for electrons and $(0.7 \pm 0.2)\%$ for muons. The scale factors are applied to the sideband region with two light leptons and an isolated track. The scale factors depend on the heavy-flavor content in the different signal regions. We parametrize this dependence as a function of the impact parameter distribution of nonisolated tracks. The $t\bar{t}$ contribution is taken from simulation.

The τ_h misidentification rate is measured in jet-dominated data by comparing the number of τ_h candidates in the signal region defined by $E_{\mathrm{cone}} < 2$ GeV to the number of nonisolated τ_h candidates, which have $6 < E_{\mathrm{cone}} < 15$ GeV. We measure the average misidentification rate as 15% with a systematic uncertainty of 30% based on the variation in different control samples. We apply this scale factor to the sideband region with two light leptons and one nonisolated τ_h candidate.

TABLE I. Observed yields for three- and four-lepton events from 19.5 fb^{-1} recorded in 2012. The channels are split by the total number of leptons (N_L), the number of τ_h candidates (N_τ), and the S_T. Expected yields are the sum of simulation and estimates of backgrounds from data in each channel. SR1–SR4 require a b-tagged jet and veto events containing Z bosons. SR5–SR8 contain events that either contain a Z boson or have no b-tagged jet. The channels are mutually exclusive. The uncertainties include statistical and systematic uncertainties. The S_T values are given in GeV.

SR	N_L	N_τ	$0 < S_T < 300$		$300 < S_T < 600$		$600 < S_T < 1000$		$1000 < S_T < 1500$		$S_T > 1500$	
			obs	exp	obs	exp	obs	exp	obs	exp	obs	exp
SR1	3	0	116	123 ± 50	130	127 ± 54	13	18.9 ± 6.7	1	1.43 ± 0.51	0	0.208 ± 0.096
SR2	3	≥ 1	710	698 ± 287	746	837 ± 423	83	97 ± 48	3	6.9 ± 3.9	0	0.73 ± 0.49
SR3	4	0	0	0.186 ± 0.074	1	0.43 ± 0.22	0	0.19 ± 0.12	0	0.037 ± 0.039	0	0.000 ± 0.03
SR4	4	≥ 1	1	0.89 ± 0.42	0	1.31 ± 0.48	0	0.39 ± 0.19	0	0.019 ± 0.026	0	0.000 ± 0.03
SR5	3	0	···	···	···	···	152	161 ± 51	15	21.0 ± 8.6	10	3.45 ± 1.77
SR6	3	1	···	···	···	···	193	150 ± 37	14	12.8 ± 3.5	0	2.04 ± 0.79
SR7	4	0	···	···	···	···	5	8.2 ± 2.6	2	0.93 ± 0.36	0	0.18 ± 0.08
SR8	4	1	···	···	···	···	2	3.2 ± 0.9	0	0.28 ± 0.13	0	0.08 ± 0.05

Another source of background leptons is internal conversions, where a virtual photon decays to a dilepton pair. These conversions produce muons almost as often as electrons, and have been discussed in detail elsewhere [20]. We measure the conversion factors of photons to light leptons in a control region (low E_T^{miss} and low hadronic activity). The ratio of the number of $\ell^+\ell^-\ell^\pm$ candidates to the number of $\ell^+\ell^-\gamma$ candidates in the Z boson decays defines the conversion factor, which is $2.1\% \pm 1.0\%$ ($0.5\% \pm 0.3\%$) for electrons (muons).

A systematic uncertainty of 4.4% in the normalization of the simulated samples accounts for imperfect knowledge of the integrated luminosity of the data sample [40]. Signal cross sections have uncertainties from 15% to 51% in stop masses between 250 GeV and 1.5 TeV, which come from the parton distribution function uncertainties and the renormalization and factorization scale uncertainties [41]. We scale the WZ and ZZ simulation samples to match data in control regions. The overall systematic uncertainty on WZ and ZZ contributions to the signal regions varies between 15% and 30% depending on the kinematics, and is the combination of the normalization uncertainties with resolution uncertainties. Muon identification efficiency uncertainty is 11% at muon p_T of 10 GeV and 0.2% at 100 GeV. For electrons the uncertainties are 14% at 10 GeV and 0.6% at 100 GeV. The uncertainty on the efficiency of the bottom-quark tagger is 6%. The uncertainty on the E_T^{miss} resolution contributes a 4% uncertainty and the jet energy scale uncertainty contributes 0.5% [42]. An uncertainty of 50% for the $t\bar{t}$ background contribution is due to the low event counts in the isolation distributions in high-S_T bins, which are used to validate the misidentification rate. We apply a 50% uncertainty to the normalization of all rare processes.

We define eight mutually exclusive signal regions (SRs) depending on the total number of leptons and the number of τ_h candidates in the event, which are defined in Table I. Since our signal does not contain any Z bosons and does contain two to four bottom quarks, in SR1–SR4, we veto events in which any opposite-sign, same-flavor dilepton pairs have an invariant mass consistent with that of the Z boson (75–105 GeV) and require at least one b-tagged jet. Each of these eight SRs is divided into five bins in S_T: [0–300], [300–600], [600–1000], [1000–1500], and [>1500] GeV. We gain additional sensitivity in regions with $S_T > 600$ GeV by removing the b-tag and Z-veto requirements for events, so the SR5–SR8 contain the events that fail one or both of these requirements.

The observed and expected yields for SR1–SR8 are shown in Table I. We also show the S_T distribution for SR1 in Fig. 1 with the background expectations from different sources shown separately. Data are in good agreement with the SM predictions everywhere. See the Supplemental Material [43] for additional S_T distributions.

We demonstrate natural SUSY with RPV couplings in a stop RPV model where the light stop decays to a top quark and intermediate on- or off-shell bino, $\tilde{t}_1 \to \tilde{\chi}_1^{0*} + t$. The bino decays to two leptons and a neutrino through the leptonic RPV interactions, $\tilde{\chi}_1^{0*} \to \ell_i + \nu_j + \ell_k$ and $\nu_i + \ell_j + \ell_k$, or through the semileptonic RPV interactions,

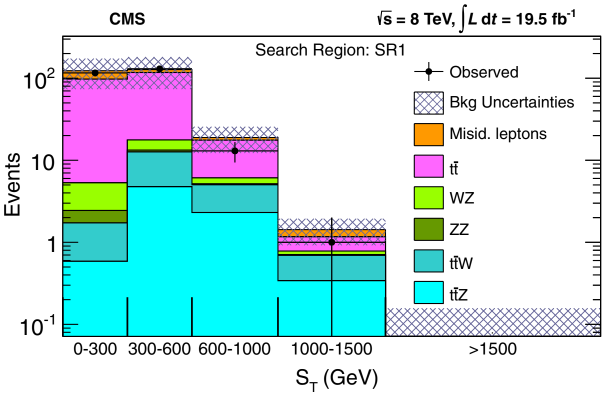

FIG. 1 (color online). The S_T distributions for SR1 including observed yields and background contributions.

$\tilde{\chi}_1^{0*} \to \ell_i + q_j + q_k$ and $\nu_i + q_j + q_k$, where the indices i, j, k refer to those appearing in Eq. (1). The stop is assumed to be right handed and RPV couplings are large enough that all decays are prompt.

We generate samples to evaluate models with simplified mass spectra and leptonic RPV couplings λ_{122} or λ_{233}. The stop masses in these samples range from 700–1250 GeV in 50 GeV steps, and bino masses range from 100–1300 GeV in 100 GeV steps. In a model with only the semileptonic RPV coupling λ'_{233}, we use stop masses 300–1000 GeV and bino masses 200–850 GeV, both in 50 GeV steps. In both cases, slepton and sneutrino masses are 200 GeV above the bino mass. Other particles are irrelevant in these models. Efficiency times acceptance figures for these models can be found in the Supplemental Material [43].

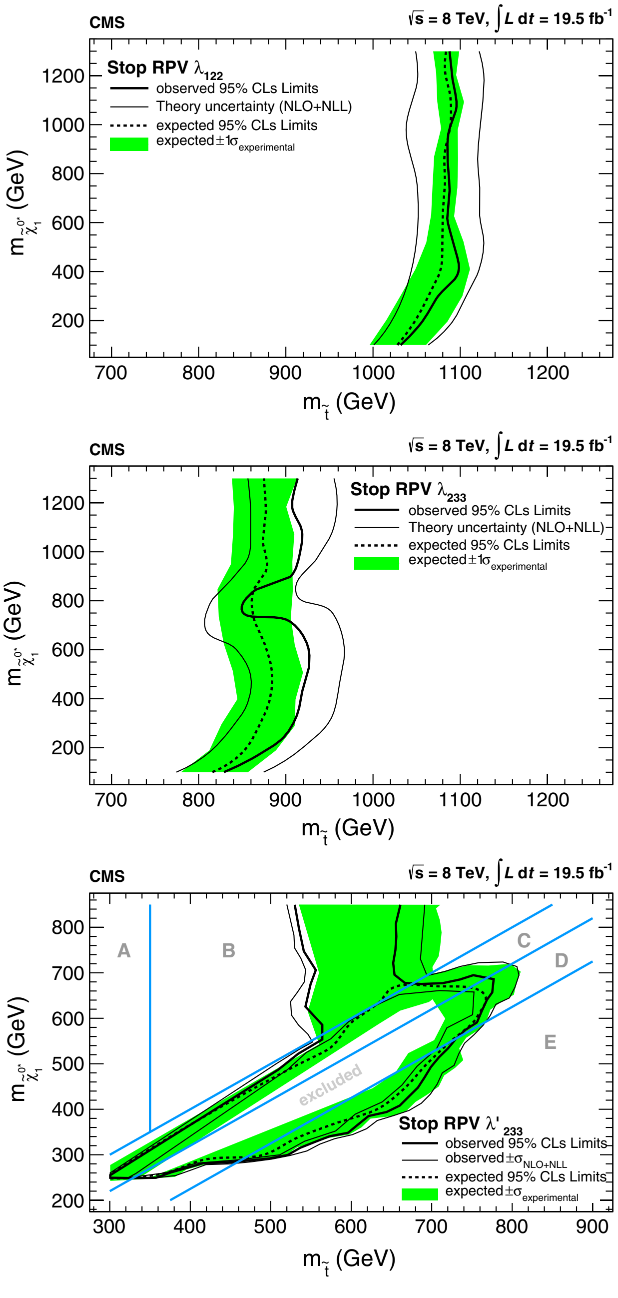

FIG. 2 (color online). The 95% confidence level limits in the stop and bino mass plane for models with RPV couplings λ_{122}, λ_{233}, and λ'_{233}. For the couplings λ_{122} and λ_{233}, the region to the left of the curve is excluded. For λ'_{233}, the region inside the curve is excluded. The different kinematic regions, A, B, C, D, and E, for the λ'_{233} exclusion are explained in Table II.

To determine which regions of phase space are excluded, we divide the channels shown in Table I by lepton flavor and perform a counting experiment using the observed event yields, the background expectations, and the signal expectations as inputs to an LHC-type CL$_s$ limit calculation [44–46]. A table with the finer binning is available in the Supplemental Material [46].

In the models with leptonic couplings, the limits are mostly independent of the bino mass, and, using the conservative minus-one standard deviation of the theoretical cross section with the observed result where the bino mass is 200 GeV, we exclude models with the stop mass below 1020 GeV when λ_{122} is nonzero, and below 820 GeV when λ_{233} is nonzero. These limits are shown in Fig. 2. There is a change in kinematics at the line $m_{\tilde{\chi}_1^0} = m_{\tilde{t}_1} - m_t$, below which the stop decay is two body, while above it is a four-body decay. Near this line, the $\tilde{\chi}_1^0$ and top are produced almost at rest, which results in soft leptons, reducing our acceptance. This loss of acceptance is more pronounced in the $\lambda_{233} \neq 0$ case and causes the loss of sensitivity near the line at $m_{\tilde{\chi}_1^0} = 800$ GeV. This feature is enhanced in the observed limit because the observed data have a larger statistical uncertainty in the relevant signal regions than the simulated signal samples.

In the semileptonic RPV model with λ'_{233}, there are several different kinematic regions, which are described in Table II. The most significant effect is when the decay $\tilde{\chi}_1^0 \to \mu + t + b$ is disfavored, reducing the number of leptons. The different regions where this effect is pronounced drive the shape of the exclusion for λ'_{233}. The area inside the curve is excluded. The observed limit is stronger than the expected one, which allows the observed exclusion region to reach into the regime where the bino decouples.

We have performed a search for RPV supersymmetry in models with top-squark pair production using a variety of

TABLE II. Kinematically allowed stop decay modes with RPV coupling λ'_{233}. The allowed neutralino decay modes for $m_t < m_{\tilde{\chi}_1^0} < m_{\tilde{t}_1}$ are $\tilde{\chi}_1^0 \to \mu t\bar{b}$ and $\nu b\bar{b}$.

Label	Kinematic region	Decay mode
A	$m_t < m_{\tilde{t}_1} < 2m_t, m_{\tilde{\chi}_1^0}$	$\tilde{t}_1 \to t\nu b\bar{b}$
B	$2m_t < m_{\tilde{t}_1} < m_{\tilde{\chi}_1^0}$	$\tilde{t}_1 \to t\mu t\bar{b}$ or $t\nu b\bar{b}$
C	$m_{\tilde{\chi}_1^0} < m_{\tilde{t}_1} < m_{W^\pm} + m_{\tilde{\chi}_1^0}$	$\tilde{t}_1 \to \ell\nu b\tilde{\chi}_1^0$ or $jjb\tilde{\chi}_1^0$
D	$m_{W^\pm} + m_{\tilde{\chi}_1^0} < m_{\tilde{t}_1} < m_t + m_{\tilde{\chi}_1^0}$	$\tilde{t}_1 \to bW^\pm\tilde{\chi}_1^0$
E	$m_t + m_{\tilde{\chi}_1^0} < m_{\tilde{t}_1}$	$\tilde{t}_1 \to t\tilde{\chi}_1^0$

multilepton final states. Good agreement between observations and SM expectations allows us to set stringent limits on the top-squark mass in models with leptonic RPV couplings λ_{122} and λ_{233}. For a bino mass of 200 GeV, these limits are 1020 GeV and 820 GeV, respectively. We also set limits in a model with the semileptonic RPV coupling λ'_{233}.

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S. Hewamanage,[131] L. M. Lebolo,[131] S. Linn,[131] P. Markowitz,[131] G. Martinez,[131] J. L. Rodriguez,[131] T. Adams,[132] A. Askew,[132] J. Bochenek,[132] J. Chen,[132] B. Diamond,[132] S. V. Gleyzer,[132] J. Haas,[132] S. Hagopian,[132] V. Hagopian,[132] K. F. Johnson,[132] H. Prosper,[132] V. Veeraraghavan,[132] M. Weinberg,[132] M. M. Baarmand,[133] B. Dorney,[133] M. Hohlmann,[133] H. Kalakhety,[133] F. Yumiceva,[133] M. R. Adams,[134] L. Apanasevich,[134] V. E. Bazterra,[134] R. R. Betts,[134] I. Bucinskaite,[134] J. Callner,[134] R. Cavanaugh,[134] O. Evdokimov,[134] L. Gauthier,[134] C. E. Gerber,[134] D. J. Hofman,[134] S. Khalatyan,[134] P. Kurt,[134] F. Lacroix,[134] D. H. Moon,[134] C. O'Brien,[134] C. Silkworth,[134] D. Strom,[134] P. Turner,[134] N. Varelas,[134] U. Akgun,[135] E. A. Albayrak,[135,yy] B. Bilki,[135,fff] W. Clarida,[135] K. Dilsiz,[135] F. Duru,[135] S. Griffiths,[135] J.-P. Merlo,[135] H. Mermerkaya,[135,ggg] A. Mestvirishvili,[135] A. Moeller,[135] J. Nachtman,[135] C. R. Newsom,[135] H. Ogul,[135] Y. Onel,[135] F. Ozok,[135,yy] S. Sen,[135] P. Tan,[135] E. Tiras,[135] J. Wetzel,[135] T. Yetkin,[135,hhh] K. Yi,[135] B. A. Barnett,[136] B. Blumenfeld,[136] S. Bolognesi,[136] D. Fehling,[136] G. Giurgiu,[136] A. V. Gritsan,[136] G. Hu,[136] P. Maksimovic,[136] M. Swartz,[136] A. Whitbeck,[136] P. Baringer,[137] A. Bean,[137] G. Benelli,[137] R. P. Kenny III,[137] M. Murray,[137] D. Noonan,[137] S. Sanders,[137] R. Stringer,[137] J. S. Wood,[137] A. F. Barfuss,[138] I. Chakaberia,[138] A. Ivanov,[138] S. Khalil,[138] M. Makouski,[138] Y. Maravin,[138] S. Shrestha,[138] I. Svintradze,[138] J. Gronberg,[139] D. Lange,[139] F. Rebassoo,[139] D. Wright,[139] A. Baden,[140] B. Calvert,[140] S. C. Eno,[140] J. A. Gomez,[140] N. J. Hadley,[140] R. G. Kellogg,[140] T. Kolberg,[140] Y. Lu,[140] M. Marionneau,[140] A. C. Mignerey,[140] K. Pedro,[140] A. Peterman,[140] A. Skuja,[140] J. Temple,[140] M. B. Tonjes,[140] S. C. Tonwar,[140] A. Apyan,[141] G. Bauer,[141] W. Busza,[141] I. A. Cali,[141] M. Chan,[141] V. Dutta,[141] G. Gomez Ceballos,[141] M. Goncharov,[141] Y. Kim,[141] M. Klute,[141] Y. S. Lai,[141] A. Levin,[141] P. D. Luckey,[141] T. Ma,[141] S. Nahn,[141] C. Paus,[141] D. Ralph,[141] C. Roland,[141] G. Roland,[141] G. S. F. Stephans,[141] F. Stöckli,[141] K. Sumorok,[141] K. Sung,[141] D. Velicanu,[141] R. Wolf,[141] B. Wyslouch,[141] M. Yang,[141] Y. Yilmaz,[141] A. S. Yoon,[141] M. Zanetti,[141] V. Zhukova,[141] B. Dahmes,[142] A. De Benedetti,[142] G. Franzoni,[142] A. Gude,[142] J. Haupt,[142] S. C. Kao,[142] K. Klapoetke,[142] Y. Kubota,[142] J. Mans,[142] N. Pastika,[142] R. Rusack,[142] M. Sasseville,[142] A. Singovsky,[142] N. Tambe,[142] J. Turkewitz,[142] L. M. Cremaldi,[143] R. Kroeger,[143] L. Perera,[143] R. Rahmat,[143] D. A. Sanders,[143] D. Summers,[143] E. Avdeeva,[144] K. Bloom,[144] S. Bose,[144] D. R. Claes,[144] A. Dominguez,[144] M. Eads,[144] R. Gonzalez Suarez,[144] J. Keller,[144] I. Kravchenko,[144] J. Lazo-Flores,[144] S. Malik,[144] F. Meier,[144] G. R. Snow,[144] J. Dolen,[145] A. Godshalk,[145] I. Iashvili,[145] S. Jain,[145] A. Kharchilava,[145] A. Kumar,[145] S. Rappoccio,[145] Z. Wan,[145] G. Alverson,[146] E. Barberis,[146] D. Baumgartel,[146] M. Chasco,[146] J. Haley,[146] A. Massironi,[146] D. Nash,[146] T. Orimoto,[146] D. Trocino,[146] D. Wood,[146] J. Zhang,[146] A. Anastassov,[147] K. A. Hahn,[147] A. Kubik,[147] L. Lusito,[147] N. Mucia,[147] N. Odell,[147] B. Pollack,[147] A. Pozdnyakov,[147] M. Schmitt,[147] S. Stoynev,[147] M. Velasco,[147] S. Won,[147] D. Berry,[148] A. Brinkerhoff,[148] K. M. Chan,[148] M. Hildreth,[148] C. Jessop,[148] D. J. Karmgard,[148] J. Kolb,[148] K. Lannon,[148] W. Luo,[148] S. Lynch,[148] N. Marinelli,[148] D. M. Morse,[148] T. Pearson,[148] M. Planer,[148] R. Ruchti,[148] J. Slaunwhite,[148] N. Valls,[148] M. Wayne,[148] M. Wolf,[148] L. Antonelli,[149] B. Bylsma,[149] L. S. Durkin,[149] C. Hill,[149] R. Hughes,[149] K. Kotov,[149] T. Y. Ling,[149] D. Puigh,[149] M. Rodenburg,[149] G. Smith,[149] C. Vuosalo,[149] G. Williams,[149] B. L. Winer,[149] H. Wolfe,[149] E. Berry,[150] P. Elmer,[150] V. Halyo,[150] P. Hebda,[150] J. Hegeman,[150] A. Hunt,[150] P. Jindal,[150] S. A. Koay,[150] D. Lopes Pegna,[150] P. Lujan,[150] D. Marlow,[150] T. Medvedeva,[150] M. Mooney,[150] J. Olsen,[150] P. Piroué,[150] X. Quan,[150] A. Raval,[150] H. Saka,[150] D. Stickland,[150] C. Tully,[150] J. S. Werner,[150] S. C. Zenz,[150] A. Zuranski,[150] E. Brownson,[151] A. Lopez,[151] H. Mendez,[151] J. E. Ramirez Vargas,[151] E. Alagoz,[152] D. Benedetti,[152] G. Bolla,[152] D. Bortoletto,[152] M. De Mattia,[152] A. Everett,[152] Z. Hu,[152] M. Jones,[152] K. Jung,[152] O. Koybasi,[152] M. Kress,[152] N. Leonardo,[152] V. Maroussov,[152] P. Merkel,[152] D. H. Miller,[152] N. Neumeister,[152] I. Shipsey,[152] D. Silvers,[152] A. Svyatkovskiy,[152] M. Vidal Marono,[152] F. Wang,[152] L. Xu,[152] H. D. Yoo,[152] J. Zablocki,[152] Y. Zheng,[152] S. Guragain,[153] N. Parashar,[153] A. Adair,[154] B. Akgun,[154] K. M. Ecklund,[154] F. J. M. Geurts,[154] W. Li,[154] B. P. Padley,[154] R. Redjimi,[154] J. Roberts,[154] J. Zabel,[154] B. Betchart,[155] A. Bodek,[155] R. Covarelli,[155] P. de Barbaro,[155] R. Demina,[155] Y. Eshaq,[155] T. Ferbel,[155] A. Garcia-Bellido,[155] P. Goldenzweig,[155] J. Han,[155] A. Harel,[155] D. C. Miner,[155] G. Petrillo,[155] D. Vishnevskiy,[155] M. Zielinski,[155] A. Bhatti,[156] R. Ciesielski,[156] L. Demortier,[156] K. Goulianos,[156] G. Lungu,[156] S. Malik,[156] C. Mesropian,[156] S. Arora,[157] A. Barker,[157] J. P. Chou,[157] C. Contreras-Campana,[157] E. Contreras-Campana,[157] D. Duggan,[157] D. Ferencek,[157] Y. Gershtein,[157] R. Gray,[157] E. Halkiadakis,[157] D. Hidas,[157] A. Lath,[157] S. Panwalkar,[157] M. Park,[157] R. Patel,[157] V. Rekovic,[157] J. Robles,[157] S. Salur,[157] S. Schnetzer,[157] C. Seitz,[157] S. Somalwar,[157] R. Stone,[157] S. Thomas,[157] P. Thomassen,[157] M. Walker,[157] G. Cerizza,[158] M. Hollingsworth,[158] K. Rose,[158] S. Spanier,[158] Z. C. Yang,[158] A. York,[158] O. Bouhali,[159,iii] R. Eusebi,[159] W. Flanagan,[159] J. Gilmore,[159] T. Kamon,[159,jjj] V. Khotilovich,[159] R. Montalvo,[159] I. Osipenkov,[159] Y. Pakhotin,[159] A. Perloff,[159] J. Roe,[159]

A. Safonov,[159] T. Sakuma,[159] I. Suarez,[159] A. Tatarinov,[159] D. Toback,[159] N. Akchurin,[160] J. Damgov,[160] C. Dragoiu,[160] P. R. Dudero,[160] C. Jeong,[160] K. Kovitanggoon,[160] S. W. Lee,[160] T. Libeiro,[160] I. Volobouev,[160] E. Appelt,[161] A. G. Delannoy,[161] S. Greene,[161] A. Gurrola,[161] W. Johns,[161] C. Maguire,[161] Y. Mao,[161] A. Melo,[161] M. Sharma,[161] P. Sheldon,[161] B. Snook,[161] S. Tuo,[161] J. Velkovska,[161] M. W. Arenton,[162] S. Boutle,[162] B. Cox,[162] B. Francis,[162] J. Goodell,[162] R. Hirosky,[162] A. Ledovskoy,[162] C. Lin,[162] C. Neu,[162] J. Wood,[162] S. Gollapinni,[163] R. Harr,[163] P. E. Karchin,[163] C. Kottachchi Kankanamge Don,[163] P. Lamichhane,[163] A. Sakharov,[163] D. A. Belknap,[164] L. Borrello,[164] D. Carlsmith,[164] M. Cepeda,[164] S. Dasu,[164] E. Friis,[164] M. Grothe,[164] R. Hall-Wilton,[164] M. Herndon,[164] A. Hervé,[164] K. Kaadze,[164] P. Klabbers,[164] J. Klukas,[164] A. Lanaro,[164] R. Loveless,[164] A. Mohapatra,[164] M. U. Mozer,[164] I. Ojalvo,[164] G. A. Pierro,[164] G. Polese,[164] I. Ross,[164] A. Savin,[164] W. H. Smith,[164] and J. Swanson[164]

(CMS Collaboration)

[1]*Yerevan Physics Institute, Yerevan, Armenia*
[2]*Institut für Hochenergiephysik der OeAW, Wien, Austria*
[3]*National Centre for Particle and High Energy Physics, Minsk, Belarus*
[4]*Universiteit Antwerpen, Antwerpen, Belgium*
[5]*Vrije Universiteit Brussel, Brussel, Belgium*
[6]*Université Libre de Bruxelles, Bruxelles, Belgium*
[7]*Ghent University, Ghent, Belgium*
[8]*Université Catholique de Louvain, Louvain-la-Neuve, Belgium*
[9]*Université de Mons, Mons, Belgium*
[10]*Centro Brasileiro de Pesquisas Fisicas, Rio de Janeiro, Brazil*
[11]*Universidade do Estado do Rio de Janeiro, Rio de Janeiro, Brazil*
[12a]*Universidade Estadual Paulista, São Paulo, Brazil*
[12b]*Universidade Federal do ABC, São Paulo, Brazil*
[13]*Institute for Nuclear Research and Nuclear Energy, Sofia, Bulgaria*
[14]*University of Sofia, Sofia, Bulgaria*
[15]*Institute of High Energy Physics, Beijing, China*
[16]*State Key Laboratory of Nuclear Physics and Technology, Peking University, Beijing, China*
[17]*Universidad de Los Andes, Bogota, Colombia*
[18]*Technical University of Split, Split, Croatia*
[19]*University of Split, Split, Croatia*
[20]*Institute Rudjer Boskovic, Zagreb, Croatia*
[21]*University of Cyprus, Nicosia, Cyprus*
[22]*Charles University, Prague, Czech Republic*
[23]*Academy of Scientific Research and Technology of the Arab Republic of Egypt, Egyptian Network of High Energy Physics, Cairo, Egypt*
[24]*National Institute of Chemical Physics and Biophysics, Tallinn, Estonia*
[25]*Department of Physics, University of Helsinki, Helsinki, Finland*
[26]*Helsinki Institute of Physics, Helsinki, Finland*
[27]*Lappeenranta University of Technology, Lappeenranta, Finland*
[28]*DSM/IRFU, CEA/Saclay, Gif-sur-Yvette, France*
[29]*Laboratoire Leprince-Ringuet, Ecole Polytechnique, IN2P3-CNRS, Palaiseau, France*
[30]*Institut Pluridisciplinaire Hubert Curien, Université de Strasbourg, Université de Haute Alsace Mulhouse, CNRS/IN2P3, Strasbourg, France*
[31]*Centre de Calcul de l'Institut National de Physique Nucleaire et de Physique des Particules, CNRS/IN2P3, Villeurbanne, France*
[32]*Université de Lyon, Université Claude Bernard Lyon 1, CNRS-IN2P3, Institut de Physique Nucléaire de Lyon, Villeurbanne, France*
[33]*Institute of High Energy Physics and Informatization, Tbilisi State University, Tbilisi, Georgia*
[34]*RWTH Aachen University, I. Physikalisches Institut, Aachen, Germany*
[35]*RWTH Aachen University, III. Physikalisches Institut A, Aachen, Germany*
[36]*RWTH Aachen University, III. Physikalisches Institut B, Aachen, Germany*
[37]*Deutsches Elektronen-Synchrotron, Hamburg, Germany*
[38]*University of Hamburg, Hamburg, Germany*
[39]*Institut für Experimentelle Kernphysik, Karlsruhe, Germany*
[40]*Institute of Nuclear and Particle Physics (INPP), NCSR Demokritos, Aghia Paraskevi, Greece*
[41]*University of Athens, Athens, Greece*
[42]*University of Ioánnina, Ioánnina, Greece*

[43]*KFKI Research Institute for Particle and Nuclear Physics, Budapest, Hungary*
[44]*Institute of Nuclear Research ATOMKI, Debrecen, Hungary*
[45]*University of Debrecen, Debrecen, Hungary*
[46]*National Institute of Science Education and Research, Bhubaneswar, India*
[47]*Panjab University, Chandigarh, India*
[48]*University of Delhi, Delhi, India*
[49]*Saha Institute of Nuclear Physics, Kolkata, India*
[50]*Bhabha Atomic Research Centre, Mumbai, India*
[51]*Tata Institute of Fundamental Research - EHEP, Mumbai, India*
[52]*Tata Institute of Fundamental Research - HECR, Mumbai, India*
[53]*Institute for Research in Fundamental Sciences (IPM), Tehran, Iran*
[54]*University College Dublin, Dublin, Ireland*
[55a]*INFN Sezione di Bari, Bari, Italy*
[55b]*Università di Bari, Bari, Italy*
[55c]*Politecnico di Bari, Bari, Italy*
[56a]*INFN Sezione di Bologna, Bologna, Italy*
[56b]*Università di Bologna, Bologna, Italy*
[57a]*INFN Sezione di Catania, Catania, Italy*
[57b]*Università di Catania, Catania, Italy*
[58a]*INFN Sezione di Firenze, Firenze, Italy*
[58b]*Università di Firenze, Firenze, Italy*
[59]*INFN Laboratori Nazionali di Frascati, Frascati, Italy*
[60a]*INFN Sezione di Genova, Genova, Italy*
[60b]*Università di Genova, Genova, Italy*
[61a]*INFN Sezione di Milano-Bicocca, Milano, Italy*
[61b]*Università di Milano-Bicocca, Milano, Italy*
[62a]*INFN Sezione di Napoli, Napoli, Italy*
[62b]*Università di Napoli "Federico II", Napoli, Italy*
[62c]*Università della Basilicata (Potenza), Napoli, Italy*
[62d]*Università G. Marconi (Roma), Napoli, Italy*
[63a]*INFN Sezione di Padova, Padova, Italy*
[63b]*Università di Padova, Padova, Italy*
[63c]*Università di Trento (Trento), Padova, Italy*
[64a]*INFN Sezione di Pavia, Pavia, Italy*
[64b]*Università di Pavia, Pavia, Italy*
[65a]*INFN Sezione di Perugia, Perugia, Italy*
[65b]*Università di Perugia, Perugia, Italy*
[66a]*INFN Sezione di Pisa, Pisa, Italy*
[66b]*Università di Pisa, Pisa, Italy*
[66c]*Scuola Normale Superiore di Pisa, Pisa, Italy*
[67a]*INFN Sezione di Roma, Roma, Italy*
[67b]*Università di Roma, Roma, Italy*
[68a]*INFN Sezione di Torino, Torino, Italy*
[68b]*Università di Torino, Torino, Italy*
[68c]*Università del Piemonte Orientale (Novara), Torino, Italy*
[69a]*INFN Sezione di Trieste, Trieste, Italy*
[69b]*Università di Trieste, Trieste, Italy*
[70]*Kangwon National University, Chunchon, Korea*
[71]*Kyungpook National University, Daegu, Korea*
[72]*Chonnam National University, Institute for Universe and Elementary Particles, Kwangju, Korea*
[73]*Korea University, Seoul, Korea*
[74]*University of Seoul, Seoul, Korea*
[75]*Sungkyunkwan University, Suwon, Korea*
[76]*Vilnius University, Vilnius, Lithuania*
[77]*Centro de Investigacion y de Estudios Avanzados del IPN, Mexico City, Mexico*
[78]*Universidad Iberoamericana, Mexico City, Mexico*
[79]*Benemerita Universidad Autonoma de Puebla, Puebla, Mexico*
[80]*Universidad Autónoma de San Luis Potosí, San Luis Potosí, Mexico*
[81]*University of Auckland, Auckland, New Zealand*
[82]*University of Canterbury, Christchurch, New Zealand*
[83]*National Centre for Physics, Quaid-I-Azam University, Islamabad, Pakistan*

[84]*National Centre for Nuclear Research, Swierk, Poland*
[85]*Institute of Experimental Physics, Faculty of Physics, University of Warsaw, Warsaw, Poland*
[86]*Laboratório de Instrumentação e Física Experimental de Partículas, Lisboa, Portugal*
[87]*Joint Institute for Nuclear Research, Dubna, Russia*
[88]*Petersburg Nuclear Physics Institute, Gatchina (St. Petersburg), Russia*
[89]*Institute for Nuclear Research, Moscow, Russia*
[90]*Institute for Theoretical and Experimental Physics, Moscow, Russia*
[91]*P. N. Lebedev Physical Institute, Moscow, Russia*
[92]*Skobeltsyn Institute of Nuclear Physics, Lomonosov Moscow State University, Moscow, Russia*
[93]*State Research Center of Russian Federation, Institute for High Energy Physics, Protvino, Russia*
[94]*University of Belgrade, Faculty of Physics and Vinca Institute of Nuclear Sciences, Belgrade, Serbia*
[95]*Centro de Investigaciones Energéticas Medioambientales y Tecnológicas (CIEMAT), Madrid, Spain*
[96]*Universidad Autónoma de Madrid, Madrid, Spain*
[97]*Universidad de Oviedo, Oviedo, Spain*
[98]*Instituto de Física de Cantabria (IFCA), CSIC-Universidad de Cantabria, Santander, Spain*
[99]*CERN, European Organization for Nuclear Research, Geneva, Switzerland*
[100]*Paul Scherrer Institut, Villigen, Switzerland*
[101]*Institute for Particle Physics, ETH Zurich, Zurich, Switzerland*
[102]*Universität Zürich, Zurich, Switzerland*
[103]*National Central University, Chung-Li, Taiwan*
[104]*National Taiwan University (NTU), Taipei, Taiwan*
[105]*Chulalongkorn University, Bangkok, Thailand*
[106]*Cukurova University, Adana, Turkey*
[107]*Middle East Technical University, Physics Department, Ankara, Turkey*
[108]*Bogazici University, Istanbul, Turkey*
[109]*Istanbul Technical University, Istanbul, Turkey*
[110]*National Scientific Center, Kharkov Institute of Physics and Technology, Kharkov, Ukraine*
[111]*University of Bristol, Bristol, United Kingdom*
[112]*Rutherford Appleton Laboratory, Didcot, United Kingdom*
[113]*Imperial College, London, United Kingdom*
[114]*Brunel University, Uxbridge, United Kingdom*
[115]*Baylor University, Waco, Texas, USA*
[116]*The University of Alabama, Tuscaloosa, Alabama, USA*
[117]*Boston University, Boston, Massachusetts, USA*
[118]*Brown University, Providence, Rhode Island, USA*
[119]*University of California, Davis, Davis, California, USA*
[120]*University of California, Los Angeles, Los Angeles, California, USA*
[121]*University of California, Riverside, Riverside, California, USA*
[122]*University of California, San Diego, La Jolla, California, USA*
[123]*University of California, Santa Barbara, Santa Barbara, California, USA*
[124]*California Institute of Technology, Pasadena, California, USA*
[125]*Carnegie Mellon University, Pittsburgh, Pennsylvania, USA*
[126]*University of Colorado at Boulder, Boulder, Colorado, USA*
[127]*Cornell University, Ithaca, New York, USA*
[128]*Fairfield University, Fairfield, Connecticut, USA*
[129]*Fermi National Accelerator Laboratory, Batavia, Illiniois, USA*
[130]*University of Florida, Gainesville, Florida, USA*
[131]*Florida International University, Miami, Florida, USA*
[132]*Florida State University, Tallahassee, Florida, USA*
[133]*Florida Institute of Technology, Melbourne, Florida, USA*
[134]*University of Illinois at Chicago (UIC), Chicago, Illinois, USA*
[135]*The University of Iowa, Iowa City, Iowa, USA*
[136]*Johns Hopkins University, Baltimore, Maryland, USA*
[137]*The University of Kansas, Lawrence, Kansas, USA*
[138]*Kansas State University, Manhattan, Kansas, USA*
[139]*Lawrence Livermore National Laboratory, Livermore, California, USA*
[140]*University of Maryland, College Park, Maryland, USA*
[141]*Massachusetts Institute of Technology, Cambridge, Massachusetts, USA*
[142]*University of Minnesota, Minneapolis, Minnesota, USA*
[143]*University of Mississippi, Oxford, Mississippi, USA*
[144]*University of Nebraska-Lincoln, Lincoln, Nebraska, USA*

[145]*State University of New York at Buffalo, Buffalo, New York, USA*
[146]*Northeastern University, Boston, Massachusetts, USA*
[147]*Northwestern University, Evanston, Indiana, USA*
[148]*University of Notre Dame, Notre Dame, Indiana, USA*
[149]*The Ohio State University, Columbus, Ohio, USA*
[150]*Princeton University, Princeton, New Jersey, USA*
[151]*University of Puerto Rico, Mayaguez, Puerto Rico*
[152]*Purdue University, West Lafayette, Indiana, USA*
[153]*Purdue University Calumet, Hammond, Indiana, USA*
[154]*Rice University, Houston, Texas, USA*
[155]*University of Rochester, Rochester, New York, USA*
[156]*The Rockefeller University, New York, New York, USA*
[157]*Rutgers, The State University of New Jersey, Piscataway, New Jersey, USA*
[158]*University of Tennessee, Knoxville, Tennessee, USA*
[159]*Texas A&M University, College Station, Texas, USA*
[160]*Texas Tech University, Lubbock, Texas, USA*
[161]*Vanderbilt University, Nashville, Tennessee, USA*
[162]*University of Virginia, Charlottesville, Virginia, USA*
[163]*Wayne State University, Detroit, Michigan, USA*
[164]*University of Wisconsin, Madison, Wisconsin, USA*

[a] Deceased.
[b] Also at Vienna University of Technology, Vienna, Austria.
[c] Also at CERN, European Organization for Nuclear Research, Geneva, Switzerland.
[d] Also at Institut Pluridisciplinaire Hubert Curien, Université de Strasbourg, Université de Haute Alsace Mulhouse, CNRS/IN2P3, Strasbourg, France.
[e] Also at National Institute of Chemical Physics and Biophysics, Tallinn, Estonia.
[f] Also at Skobeltsyn Institute of Nuclear Physics, Lomonosov Moscow State University, Moscow, Russia.
[g] Also at Universidade Estadual de Campinas, Campinas, Brazil.
[h] Also at California Institute of Technology, Pasadena, CA, USA.
[i] Also at Laboratoire Leprince-Ringuet, Ecole Polytechnique, IN2P3-CNRS, Palaiseau, France.
[j] Also at Zewail City of Science and Technology, Zewail, Egypt.
[k] Also at Suez Canal University, Suez, Egypt.
[l] Also at Cairo University, Cairo, Egypt.
[m] Also at Fayoum University, El-Fayoum, Egypt.
[n] Also at British University in Egypt, Cairo, Egypt.
[o] Now at Ain Shams University, Cairo, Egypt.
[p] Also at National Centre for Nuclear Research, Swierk, Poland.
[q] Also at Université de Haute Alsace, Mulhouse, France.
[r] Also at Joint Institute for Nuclear Research, Dubna, Russia.
[s] Also at Brandenburg University of Technology, Cottbus, Germany.
[t] Also at The University of Kansas, Lawrence, Kansas, USA.
[u] Also at Institute of Nuclear Research ATOMKI, Debrecen, Hungary.
[v] Also at Eötvös Loránd University, Budapest, Hungary.
[w] Also at Tata Institute of Fundamental Research - EHEP, Mumbai, India.
[x] Also at Tata Institute of Fundamental Research - HECR, Mumbai, India.
[y] Now at King Abdulaziz University, Jeddah, Saudi Arabia.
[z] Also at University of Visva-Bharati, Santiniketan, India.
[aa] Also at University of Ruhuna, Matara, Sri Lanka.
[bb] Also at Isfahan University of Technology, Isfahan, Iran.
[cc] Also at Sharif University of Technology, Tehran, Iran.
[dd] Also at Plasma Physics Research Center, Science and Research Branch, Islamic Azad University, Tehran, Iran.
[ee] Also at Università degli Studi di Siena, Siena, Italy.
[ff] Also at Universidad Michoacana de San Nicolas de Hidalgo, Morelia, Mexico.
[gg] Also at Faculty of Physics, University of Belgrade, Belgrade, Serbia.
[hh] Also at Facoltà Ingegneria, Università di Roma, Roma, Italy.
[ii] Also at Scuola Normale e Sezione dell'INFN, Pisa, Italy.

[jj]Also at INFN Sezione di Roma, Roma, Italy.
[kk]Also at University of Athens, Athens, Greece.
[ll]Also at Rutherford Appleton Laboratory, Didcot, United Kingdom.
[mm]Also at Paul Scherrer Institut, Villigen, Switzerland.
[nn]Also at Institute for Theoretical and Experimental Physics, Moscow, Russia.
[oo]Also at Albert Einstein Center for Fundamental Physics, Bern, Switzerland.
[pp]Also at Gaziosmanpasa University, Tokat, Turkey.
[qq]Also at Adiyaman University, Adiyaman, Turkey.
[rr]Also at Cag University, Mersin, Turkey.
[ss]Also at Mersin University, Mersin, Turkey.
[tt]Also at Izmir Institute of Technology, Izmir, Turkey.
[uu]Also at Ozyegin University, Istanbul, Turkey.
[vv]Also at Kafkas University, Kars, Turkey.
[ww]Also at Suleyman Demirel University, Isparta, Turkey.
[xx]Also at Ege University, Izmir, Turkey.
[yy]Also at Mimar Sinan University, Istanbul, Istanbul, Turkey.
[zz]Also at Kahramanmaras Sütcü Imam University, Kahramanmaras, Turkey.
[aaa]Also at School of Physics and Astronomy, University of Southampton, Southampton, United Kingdom.
[bbb]Also at INFN Sezione di Perugia, Università di Perugia, Perugia, Italy.
[ccc]Also at Utah Valley University, Orem, UT, USA.
[ddd]Also at Institute for Nuclear Research, Moscow, Russia.
[eee]Also at University of Belgrade, Faculty of Physics and Vinca Institute of Nuclear Sciences, Belgrade, Serbia.
[fff]Also at Argonne National Laboratory, Argonne, IL, USA.
[ggg]Also at Erzincan University, Erzincan, Turkey.
[hhh]Also at Yildiz Technical University, Istanbul, Turkey.
[iii]Also at Texas A&M University at Qatar, Doha, Qatar.
[jjj]Also at Kyungpook National University, Daegu, Korea.